PHYSICAL PROPERTIES OF SPECTROSCOPICALLY-CONFIRMED GALAXIES AT $Z \ge 6$. II. MORPHOLOGY OF THE REST-FRAME UV CONTINUUM AND LY- α EMISSION^{$\overline{*}$}

Linhua Jiang^{1,2,9}, Eiichi Egami², Xiaohui Fan², Rogier A. Windhorst¹, Seth H. Cohen¹, Romeel Davé^{2,3}, Kristian Finlator^{4,10}, Nobunari Kashikawa⁵, Matthew Mechtley¹, Masami Ouchi^{6,7}, and Kazuhiro Shimasaku⁸

Draft version January 5, 2018

ABSTRACT

We present a detailed structural and morphological study of a large sample of spectroscopicallyconfirmed galaxies at $z \ge 6$, using deep HST near-IR broad-band images and Subaru optical narrowband images. The galaxy sample consists of 51 Ly α emitters (LAEs) at $z \simeq 5.7$, 6.5, and 7.0, and 16 Lyman-break galaxies (LBGs) at $5.9 \le z \le 6.5$. These galaxies exhibit a wide range of restframe UV continuum morphology in the HST images, from compact features to multiple component systems. The fraction of merging/interacting galaxies reaches $40\% \sim 50\%$ at the brightest end of $M_{1500} \leq -20.5$ mag. The intrinsic half-light radii $r_{\rm hl,in}$, after correction for PSF broadening, are roughly between $r_{\rm hl,in} \simeq 0.000$ (0.3 kpc) and 0.0000 (1.7 kpc) at $M_{1500} \leq -19.5$ mag. The median $r_{\rm hl,in}$ value is 0".16 (~0.9 kpc). This is consistent with the sizes of bright LAEs and LBGs at $z \ge 6$ in previous studies. In addition, more luminous galaxies tend to have larger sizes, exhibiting a weak size-luminosity relation $r_{\rm hl,in} \propto L^{0.14}$ at $M_{1500} \leq -19.5$ mag. The slope of 0.14 is significantly flatter than those in fainter LBG samples. We discuss the morphology of $z \ge 6$ galaxies with nonparametric methods, including the CAS system and the Gini and M_{20} parameters, and demonstrate their validity through simulations. We search for extended Ly α emission halos around LAEs at $z \simeq 5.7$ and 6.5, by stacking a number of narrow-band images. We do not find evidence of extended halos predicted by cosmological simulations. Such $Ly\alpha$ halos, if they exist, could be weaker than predicted. Finally, we investigate any positional misalignment between UV continuum and $Ly\alpha$ emission in LAEs. While the two positions are generally consistent, several merging galaxies show significant positional differences. This is likely caused by a disturbed ISM distribution due to merging activity.

Subject headings: cosmology: observations — galaxies: evolution — galaxies: high-redshift

1. INTRODUCTION

Galaxy structural and morphological studies provide basic apparent information about galaxies. Nearby galaxies are generally classified into three broad categories: spiral, elliptical, and irregular. The majority of

² Steward Observatory, University of Arizona, 933 North Cherry Avenue, Tucson, AZ 85721, USA

³ Physics Department, University of the Western Cape, 7535 Bellville, Cape Town, South Africa

Dark Cosmology Centre, Niels Bohr Institute, University of Copenhagen, DK-2100 Copenhagen Ø, Denmark

⁵ Optical and Infrared Astronomy Division, National Astronomical Observatory, Mitaka, Tokyo 181-8588, Japan ⁶ Institute for Cosmic Ray Research, The University of Tokyo,

5-1-5 Kashiwanoha, Kashiwa, Chiba 277-8582, Japan 7 Kavli Institute for the Physics and Mathematics of the Universe, The University of Tokyo, 5-1-5 Kashiwanoha, Kashiwa, Chiba 277-8583, Japan

⁸ Department of Astronomy, University of Tokyo, Hongo, Tokyo 113-0033, Japan

Hubble Fellow.

¹⁰ DARK fellow.

luminous nearby galaxies ($z \leq 0.1$) are spirals and ellipticals (Abraham & van den Bergh 2001). At higher redshift, galaxies are less well developed, and the fraction of irregular galaxies increases steadily (e.g., Driver et al. 1995, 1998). In the redshift range of 0.1 < z < 1, galaxy morphology and structure have also been well studied (e.g., Brinchmann et al. 1998; Lilly et al. 1998; Schade et al. 1999; Carlberg et al. 2000; Le Fèvre et al. 2000; van den Bergh et al. 2000, and references therein). These galaxies show more disturbed structures than nearby galaxies do in the rest-frame UV and optical (e.g., Abraham & van den Bergh 2001; Windhorst et al. 2002; Taylor-Mager et al. 2007; Blanton & Moustakas 2009; Shi et al. 2009). The fraction of irregular galaxies increases from less than 10% at $z \leq 0.5$ to $\sim 30\%$ at $z \simeq 1$ (e.g., Brinchmann et al. 1998; van den Bergh et al. 2000). In addition, more galaxies were identified as merging systems, reflecting the hierarchical build-up of galaxies and mass assembly in the cold dark matter (CDM) scenario (White & Rees 1978; Cole et al. 2000). For example, nearly 20% of $z \simeq 1$ galaxies in the Le Fèvre et al. (2000) sample are in close pairs.

For galaxies at $z \geq 2-3$, morphological classification is challenging, as most galaxies appear peculiar. Galaxies are also smaller towards higher redshift (Ferguson et al. 2004). In addition, galaxies appear much fainter due to the cosmological $(1 + z)^4$ surface brightness dimming. Traditional classifications, including Hubble's tuningfork system, are no longer practical at these higher redshifts. Therefore, nonparametric methods such as

Based in part on observations made with the NASA/ESA Hubble Space Telescope, obtained from the data archive at the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Inc. under NASA contract NAS 5-26555. Based in part on observations made with the Spitzer Space Telescope, which is operated by the Jet Propulsion Laboratory, California Institute of Technology under a contract with NASA. Based in part on data collected at Subaru Telescope and obtained from the SMOKA, which is operated by the Astronomy Data Center, National Astronomical Observatory of Japan.

School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287-1504, USA; linhua.jiang@asu.edu

the CAS system (Conselice 2003) and the Gini and M_{20} parameters (Lotz et al. 2004) play an important Most morphological and structural analyses role. in this redshift range were done in the GOODS fields (Giavalisco et al. 2004), because of the high-quality Hubble Space Telescope (HST) data (e.g., Lowenthal et al. 1997; Ravindranath et al. 2006; Law et al. 2007; Cassata et al. 2010). In the rest-frame UV, $z \geq 3$ galaxies are usually compact (from one to several kpc), but many of them display extended features or multiple clumps in deep HST images (e.g., Giavalisco et al. 2005; Ravindranath et al. 1996:Venemans et al. 2006; Pirzkal et al. 2007; Conselice & Arnold 2009; Cooke et al. 2010; Gronwall et al. 2011; Law et al. 2012).For example, the Ravindranath et al. (2006)sample contains thousands of photometrically-selected LBGs at $z \ge 2.5$, and about 30% of them have In Law et al. (2012), a sample of multiple cores. spectroscopically-confirmed galaxies in a similar redshift range also showed a high fraction of interacting systems.

In the highest-redshift range $z \ge 6$, morphological studies become even more difficult. Galaxies appear very faint, and their rest-frame UV light moves to the near-IR wavelength range, where telescope resolution is poorer. A typical galaxy occupies only a few pixels even in HST near-IR images, so size is usually the only physical parameter that can be reliably measured in the literature. Studies based on photometrically-selected galaxies have shown that $z \ge 6$ galaxies are generally very compact, and most of them are just barely spatially resolved. For example, Oesch et al. (2010) reported the sizes of 16 LBGs at $z \ge 7$ in the Hubble Ultra-Deep Field (HUDF). They found that only two in their sample show extended features, and the rest are very compact (< 1 kpc). This sample is very faint. Observations of a handful brighter galaxies with spectroscopic redshifts also suggest compact morphology, with a typical size of < 1 kpc (e.g., Stanway et al. 2004; Dow-Hygelund et al. 2007; Cowie et al. 2011). Note that galaxy size is correlated with physical properties, such as mass and luminosity at low redshift. Such relations may still exist in high-redshift galaxies, but could have evolved with time (e.g., Grazian et al. 2012; Mosleh et al. 2012).

In this paper, we will carry out a structural and morphological study of a sample of 67 galaxies at $z \ge 6$. The sample is the largest collection of spectroscopicallyconfirmed galaxies in this redshift range, including 51 $Ly\alpha$ emitters (LAEs) and 16 Lyman-break galaxies (LBGs). This paper is the second in a series presenting the physical properties of these galaxies. In the first paper of the series (Jiang et al. 2013, hereafter Paper I), we presented deep Subaru optical and HST near-IR data. We also derived various rest-frame UV continuum and $Ly\alpha$ emission properties, including UV-continuum slope β , the Ly α rest-frame equivalent width (EW), and star formation rates (SFRs). These galaxies have steep UV continuum slopes roughly between $\beta \simeq -1.5$ and -3.5, with a mean value of $\beta \simeq -2.3$. They have a range of $Ly\alpha$ EW from ~10 to ~200 Å. Their SFRs are moderate from a few to a few tens solar masses per year. In this paper, we will study the structure and morphology of their rest-frame UV-continuum emission based on our HST images, and of their $Ly\alpha$ emission based on our

ground-based narrow-band images.

The layout of the paper is as follows. In Section 2, we briefly review our galaxy sample and the optical and near-IR data that will be used for the paper. We measure the structure and morphology of UV continuum emission in Section 3, and of the Ly α emission in Section 4. We then discuss our results and summarize the paper in Section 5. Throughout the paper we use a Λ -dominated flat cosmology with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.3$, and $\Omega_{\Lambda} = 0.7$ (Komatsu et al. 2011). All magnitudes are on the AB system (Oke & Gunn 1983).

2. GALAXY SAMPLE AND DATA

In Section 2 of Paper I, we provided a detailed description of our galaxy sample and the multi-wavelength data that we studied. Here we summarize the information below. There are a total of 67 spectroscopically-confirmed galaxies in our sample: 62 are from the Subaru Deep Field (SDF; Kashikawa et al. 2004), and the remaining 5 are from the Subaru XMM-Newton Deep Survey field (SXDS; Furusawa et al. 2008). They represent the most luminous galaxies in terms of $Lv\alpha$ luminosity (for LAEs) or UV continuum luminosity (for LBGs) in this redshift range. The SDF galaxy sample contains 22 LAEs at $z \simeq$ 5.7 (Shimasaku et al. 2006; Kashikawa et al. 2011), 25 LAEs at $z \simeq 6.5$ (Taniguchi et al. 2005; Kashikawa et al. 2006, 2011), and a LAE at z = 6.96 (Ive et al. 2006). The LAEs at $z \simeq 5.7$ and 6.5 have a relatively uniform magnitude limit of 26 mag in the narrow bands NB816 and NB921, and thus make a well-defined sample. The SDF sample also contains 14 LBGs at 5.9 \leq $z \leq 6.5$ (Nagao et al. 2004, 2005, 2007; Ota et al. 2008; Jiang et al. 2011; Toshikawa et al. 2012). The LBG candidates in these studies were selected with different criteria, and have a rather inhomogeneous depth. The SXDS sample consists of five galaxies, including two LBGs at $z \simeq 6$ (Curtis-Lake et al. 2012) and three LAEs at $z \simeq 6.5$ (Ouchi et al. 2010). All these galaxies are listed in Table 1 of Paper I.

The SDF and SXDS were observed with Subaru Suprime-Cam (Kashikawa et al. 2004; Furusawa et al. 2008). They have extremely deep optical images in a series of broad and narrow bands. Public stacked images are available for the two fields, but the public data do not include the images taken recently. In Paper I, we produced our own stacked images in six broad bands (BVRi'z'y) and three narrow bands (NB816, NB921, and NB973) by including all available data in the archive. Our stacked images have great depth with excellent PSF full width at half-maxima (FWHMs) of 0.5' - 0.7'. The near-IR imaging data for the SDF galaxies are from three HST GO programs 11149 (PI: E. Egami), 12329 and 12616 (PI: L. Jiang). The HST observations were made with a mix of instruments and depth. The majority of the galaxies were observed with WFC3 in the F125W (hereafter J_{125}) and F160W (hereafter H or H_{160}) bands. The typical integration time was two HST orbits (roughly 5400 s) per band. This provides a depth of ~ 27.5 mag $(5\sigma \text{ detection})$ in the J_{125} band and ~ 27.1 mag in the H_{160} band (see also Windhorst et al. 2011). The pixel size in the final reduced WFC3 images is 0".06. Several SDF galaxies were observed with NICMOS in the F110W (hereafter J_{110}) and H_{160} bands. The typical integration time was also two HST orbits. The depth in

the two bands are ~ 26.4 mag and ~ 26.1 mag, respectively. The pixel size in the final reduced NICMOS images is 0".1. The five SXDS galaxies were covered by the UKIDSS Ultra-Deep Survey (UDS). Their *HST* WFC3 near-IR data were obtained from the Cosmic Assembly Near-infrared Deep Extragalactic Legacy Survey (CAN-DELS; Grogin et al. 2011; Koekemoer et al. 2011). The exposure depth of the CANDELS UDS data is 1900 s in the J_{125} band and 3300 s in the H_{160} band.

The majority of the galaxies in our sample were detected with high significance in the near-IR images. Only 15 of them — among the faintest in the optical — have weak detections ($< 5\sigma$) in the *J* band (J_{125} or J_{110}). In Table 1 of Paper I, we listed the optical and near-IR photometry of the galaxies. In Table 2 of Paper I, we presented basic physical properties, including the restframe UV continuum luminosity and slope β , the Ly α luminosity and EW, and SFR, etc. The thumbnail images of all the galaxies are provided in Appendix A of Paper I.

3. UV CONTINUUM MORPHOLOGY

In this section, we will derive structural and morphological parameters for the galaxies in our sample. Although our galaxies represent the most luminous galaxies at z > 6, they appear faint and small compared to lowerredshift galaxies. The majority of them are point-like sources in the Subaru optical images. Even in the HSTWFC3 images they usually occupy a very limited number of pixels. Hence, the study of these distant galaxies is challenging. In previous literature, galaxy size was often the only parameter that could be reliably measured for z > 6 galaxies. Classifications for nearby galaxies, such as the classical Hubble's tuning-fork system, cannot be applied to these objects. In this section, we will measure the sizes of our galaxies, and try to characterize their morphology using nonparametric methods, such as the CAS system (Conselice 2003) and the Gini and M_{20} parameters (Lotz et al. 2004). These methods are primarily used for low-redshift galaxies, though they have already been used for galaxies at $z = 4 \sim 6$ (e.g., Pirzkal et al. 2007; Conselice & Arnold 2009). We will also study interacting/merging systems in our sample.

In order to calculate the above parameters, we took advantage of all our HST images. For each galaxy in our sample, we combined (the weighted average) its Jand H-band images and made a stacked HST image to improve the signal-to-noise (S/N) ratio. By doing this we assume that the effect of the morphological k-correction the dependence of galaxy structure on wavelength – is negligible in the wavelength range considered. This is because the J and H bands cover a similar restframe UV wavelength range (~ 1780 Å vs. ~ 2200 Å) for $z \simeq 6$ galaxies. Our further analyses were then based on the stacked images. The J and H bands do not cover $Ly\alpha$ emission for our galaxies, so the morphology in the stacked images is purely from their UV continuum emission (other nebular lines can be safely ignored in general; see also Cai et al. 2011 and Kashikawa et al. 2012). In Figure 1, we show the thumbnail images of 44 (out of 67) galaxies that have more than 10σ detections of their total fluxes in the stacked images. We will focus on these 44 galaxies in this section. Note that we excluded object no. 12, since it overlaps with a bright foreground star,

as explained in Paper I.

We do not show in Figure 1 the galaxies with $< 10\sigma$ detections. Most of them are very faint. The others have shallower images (e.g., 1-orbit depth). We would not obtain reliable morphological information for these individual faint galaxies, so we combined their images and made a single stacked image. The individual images were scaled before stacking, so that the galaxies all have the same magnitude. The final stacked image has a much higher S/N ratio.

3.1. Size

We use half-light radius $r_{\rm hl}$ to describe the size of a galaxy. The half-light radius $r_{\rm hl}$ is a radius enclosing a half of the total light. In Figure 2, the upper panel shows the measured $r_{\rm hl}$ as the function of M_{1500} , the absolute AB magnitude of the continuum at rest-frame 1500 Å. The measured $r_{\rm hl}$ were calculated from elliptical apertures using SExtractor (Bertin & Arnouts 1996). Physical quantities such as M_{1500} were derived in Paper I. The blue and red circles in Figure 2 represent the LAEs at $z \simeq 5.7$ and 6.5 (including $z \simeq 7$), respectively, and the green circles represent the LBGs at $z \simeq 6$. We do not include object no. 67, since it has three well separated components and its $r_{\rm hl}$ is likely meaningless. The black star indicates the stacked image of the faint galaxies described above. The measurement uncertainties in the upper panel were derived from the simulations described below. These galaxies roughly span a luminosity range of $-22 \leq M_{1500} \leq -19.5$, and a radius range of $0'.1 \leq r_{\rm hl} \leq 0''.3$ (or 0.6–1.7 kpc) without correction for PSF broadening. The median $r_{\rm hl}$ value is 0'.19 (~ 1.1 kpc). Note that the sample in Figure 2 is dominated by the galaxies at $M_{1500} \leq -20$ mag. The stacked image of the faint galaxies has a relatively smaller radius $r_{\rm hl} \simeq 0^{\prime\prime}.14 \ (\sim 0.8 \ {\rm kpc})$. The median $r_{\rm hl}$ value for the whole sample of 67 galaxies is $0^{\prime\prime}_{...}16 \ (\sim 0.9 \text{ kpc})$, if we assume that the faint galaxies with $M_{1500} \ge -19.5$ mag in this sample have $r_{\rm hl}$ smaller than this median value.

We estimated measurement uncertainties for $r_{\rm hl}$ from simulations. For each galaxy, we made a model galaxy using the two-dimensional fitting algorithm GALFIT (Peng et al. 2002). A single component with a Sérsic function was fitted to the galaxy. This noiseless mock galaxy was put back at 300 ~ 400 random positions (one position at a time) in the blank regions of the same science image. We then measured $r_{\rm hl}$ using the same way as we did for our real $z \ge 6$ galaxies. We denoted the standard deviation of the measured $r_{\rm hl}$ as the measurement uncertainty of $r_{\rm hl}$ for this galaxy. The measurement uncertainties of $r_{\rm hl}$ in our sample (shown as the error bars in the upper panel of Figure 2) have a median value of 10%.

We estimated intrinsic half-light radius $r_{\rm hl,in}$ and systematic uncertainties from simulations as well. Because of the PSF broadening and lower surface brightness at larger sizes, the observed (or measured) $r_{\rm hl}$ is a complex function of galaxy size, brightness, and intrinsic profile. Simulations have been widely used to investigate these systematic effects (e.g. Driver et al. 2005; Häussler et al. 2007; Cibinel et al. 2012; Grazian et al. 2012; van der Wel et al. 2012). We started with the GALFIT parameters obtained above, and considered four parameters here, including the Sérsic index n, the axis

ratio b/a, brightness, and $r_{\rm hl,in}$. Based on these parameters, we produced a large set of mock galaxies in a grid of $n, b/a, M_{1500}$, and $r_{hl,in}$. The values of n were chosen to be 1, 1.5, and 2 times the measured n from our galaxies, and the values of b/a were chosen to be 1, 2/3, 1/3times the measured b/a from our galaxies. This is because the PSF broadening could largely decrease n and increase b/a in low-resolution images. The luminosity coverage $(-22.4 \leq M_{1500} \leq -19.4)$ and $r_{\rm hl,in}$ coverage $(0''.03 \le r_{\rm hl,in} \le 0''.4)$ that we chose are roughly consistent with the actual coverage of our galaxy sample. The mock galaxies were oversampled so that their $r_{\rm hl,in}$ sizes were at least 20 pixels. Then they were convolved with PSF images and rebinned to match the pixel scales of our HST images. Finally, each of the rebinned mock galaxies was placed at many (> 300) random positions in the blank regions of our science images. The $r_{\rm hl}$ of this mock galaxy is the median $r_{\rm hl}$ value measured at these random positions.

Figure 3 shows part of our simulation results. It illustrates the measured $r_{\rm hl}$ as a function of $r_{\rm hl,in}$ at four different magnitudes. It clearly shows that at small sizes, $r_{\rm hl}$ is significantly larger than $r_{\rm hl,in}$ due to the PSF broadening. At large sizes, however, $r_{\rm hl}$ starts to fall short of $r_{\rm hl,in}$, because we start to loose low surface brightness pixels at large sizes by detection. This happens at smaller sizes for fainter galaxies. At the faintest magnitude $M_{1500} = -19.5 \text{ mag}$, $r_{\rm hl}$ will never exceed 0."3 in our images, regardless of $r_{\rm hl,in}$. On the other hand, if faint galaxies are always small, as seen in deeper HSTimages (e.g., Windhorst et al. 2008; Oesch et al. 2010; Grazian et al. 2012), their $r_{\rm hl}$ should not be significantly smaller than $r_{\rm hl,in}$ (e.g. $r \leq 0^{\prime\prime}2$ at $M_{1500} = -19.5$ mag in Figure 3). Note that Figure 3 is similar to Figure 4 in Grazian et al. (2012). Another simple way to correct for PSF broadening is to estimate $r_{\rm hl,in}$ in quadrature, i.e., $r_{\rm hl,in}^2 = r_{\rm hl}^2 - r_{\rm PSF}^2$, where $r_{\rm PSF}$ is the PSF radius. Figure 3 shows that this equation is a good approximation at small sizes and/or high luminosities (green dashed lines), but it underestimates $r_{\rm hl,in}$ elsewhere.

The galaxy sizes in our sample roughly agree with those of high-redshift LAEs and LBGs with similar luminosities in the previous literature. For example, Pirzkal et al. (2007) found that the average $r_{\rm hl}$ for a sample of luminous LAEs at $z \sim 5$ is 0".17. Taniguchi et al. (2009) found a median $r_{\rm hl}$ of 0".15 for LAEs at $z \sim 5.7$. In the Hathi et al. (2008) and Conselice & Arnold (2009) LBG samples of $z = 4 \sim 6$ galaxies, the $r_{\rm hl}$ ranges are 0".1 \sim 0".3, similar to the $r_{\rm hl}$ range in our sample. Previous studies have shown that the galaxy size roughly scales with redshift as $(1+z)^{-m}$, with *m* close to 1.1–1.2 (e.g. Ferguson et al. 2004; Bouwens et al. 2006; Oesch et al. 2010; Mosleh et al. 2012), so the size of galaxies evolves slowly at high redshift. This is the reason that high-redshift galaxies have a similar size range. Malhotra et al. (2012) found, however, that LAEs have a roughly constant size in the redshift range of 2.25 < z < 6, and do not show a size-redshift relation. While our sample does not have a large redshift coverage, our galaxy sizes are well consistent with those in their sample.

3.1.1. Size-luminosity relation

Figure 2 shows that brighter objects appear to have larger $r_{\rm hl}$, meaning that more luminous galaxies tend to have larger physical sizes. This size-luminosity relation has been reported for both low- and high-redshift starforming galaxies (e.g., Taniguchi et al. 2009; Oesch et al. 2010; Grazian et al. 2012; Ono et al. 2012). For example, with a large sample photometrically-selected LBGs at $z \sim 7$ in the CANDELS fields, Grazian et al. (2012) found a strong relation $r_{\rm hl} \propto L^{\alpha}$, with slope $\alpha \simeq 1/2$. In Figure 2, we illustrate the size-luminosity relation by displaying the best log-linear fits (dashed lines). The best fitting results in the two panels are $r_{\rm hl} \propto L^{0.11\pm0.02}$ and $r_{\rm hl,in} \propto L^{0.14\pm0.03}$, respectively. Our slopes are much flatter than that in Grazian et al. (2012) and those in other fainter LBG samples. The reason is that our galaxy sample is much brighter. Our relation is derived from galaxies in the luminosity range of $M_{1500} \leq -19.5$ mag, while the Grazian et al. (2012) sample covers a range of $M_{\rm UV} \leq -18$ mag. Their relation largely depends on the galaxies fainter than -19.5 mag, as seen in Figure 9 of their paper. In the brighter galaxies, their $r_{\rm hl}$ (or $r_{\rm hl,in}$) shows less of a trend with luminosity, as also pointed out by Grazian et al. (2012). In fact, our best fit to the galaxies with $M_{1500} \leq -20 \text{ mag}$ (dash-dotted line in Figure 2) gives a nearly flat slope of $\alpha = -0.06 \pm 0.03$, suggesting little correlation between size and luminosity in the most luminous galaxies.

It should be pointed out that the relation between the measured size and luminosity could be affected by systematic effects shown in Figure 3. Figure 4 illustrates how such effects shape the $r_{\rm hl}$ - M_{1500} relation for mock galaxies from the simulations above (see also e.g., Cibinel et al. 2012). The open circles indicate the intrinsic sizes $r_{\rm hl,in}$, and the filled circles are the measured sizes $r_{\rm hl}$ at different luminosities. The two colors red and blue indicate two different HST images that the mock galaxies are placed in. The two panels are for two sets of Sérsic index n: 1 and 2 times the measured n from our $z \geq 6$ galaxies. Figure 4 shows that fainter galaxies (with the same intrinsic size) appear to be smaller, as already shown in Figure 3. For the same reason, larger galaxies (with the same intrinsic luminosity) appear to be slightly fainter. We have taken into account these effects in Figure 2. These effects do not have significant impact on the size-luminosity relation in Figure 2: the corrected relation does not become flatter, because this relation also depends on other factors such as the source distribution.

Finally, we point out that the size-luminosity relation in our sample is not affected by a possible selection effect, i.e., that we may have missed some faint galaxies with large sizes during galaxy candidate selection. The reason is that these galaxies were selected in optical images. They are bright and point-like sources in the optical broad bands (for LBGs) or narrow bands (for LAEs). Also note that the exclusion of faint galaxies in our sample does not introduce bias to our results. Our *HST* near-IR data have relatively uniform depth (two orbits per band per pointing), so the 10σ cut indeed puts a flux limit on M_{1500} in Figure 2, which does not affect our results in the bright region of $M_{1500} < -19.5$ mag.

3.2. Nonparametric Measurements of Morphology

In this subsection, we will characterize galaxy structure and morphology using nonparametric methods, including the CAS (Concentration, Asymmetry, and Smoothness) system (Conselice 2003), the Gini coefficient G, and the M_{20} parameters (Lotz et al. 2004). These methods have been widely used for lower-redshift galaxies. They usually provide reliable description of galaxy structure, and are able to distinguish different types of galaxies. To obtain accurate measurements of these quantities, two criteria are often required: high S/N ratios and large object sizes compared to the PSF size. For low-redshift galaxies, especially those in HST images, the sizes of galaxies are many times larger than PSF, so the two criteria are naturally met in deep HST images. At higher redshift, galaxies are fainter and smaller, so it is difficult to meet the two criteria, and these parameters become less reliable. We investigate the reliability of these measurements for our sample in detail below.

3.2.1. CAS, Gini, and M₂₀ parameters

Concentration (C) measures how compact the galaxy light profile is. We adopted the commonly used definition $C = 5 \log (r_{80}/r_{20})$ (e.g., Bershady et al. 2000; Conselice 2003), where r_{80} and r_{20} are the radii that contain 80% and 20% of the total galaxy flux, respectively. Asymmetry (A) measures how rotationally symmetric a galaxy is (Abraham et al. 1996; Conselice 2003). It is calculated by subtracting the image rotated by 180° from the original galaxy image. Smoothness (S) or clumpiness measures how clumpy a galaxy is (Conselice 2003). Conselice & Arnold (2009) found that S fails to well describe clumpiness for $z = 4 \sim 6$ galaxies, so we did not calculate S in this paper. The M_{20} parameter, or the second-order moment of the brightest 20% of the galaxy, is similar to the concentration \tilde{C} , and measures how the galaxy light is concentrated (Lotz et al. 2004). The Gini coefficient (G) describes how even the galaxy light distribution is (Abraham et al. 2003; Lotz et al. 2004). We computed G and M_{20} as described by Lotz et al. (2004).

The measurements of these parameters are shown in Figure 5. The measurement uncertainties were estimated from simulations using the same method as we did for the uncertainties of $r_{\rm hl}$. We took the model galaxies obtained in Section 3.1, and put them at many random positions in the blank regions of our HST images. We then measured $CAGM_{20}$ at each position, and calculated the standard deviations of these parameters. The standard deviations, or measurement uncertainties, are shown as the error bars in Figure 5. These uncertainties include the effects of S/N in the images, but do not account for

systematic uncertainties associated with sparse spatial sampling of distant small and faint sources which we investigate in Section 3.2.2.

Compared to low-redshift galaxies, our galaxies are located in a much narrower range in the parameter space (Figure 5; see also Figures 6 and 7). They appear to be less concentrated and more asymmetric, and their light distribution is more even. This is likely because the measured quantities have been substantially affected by the low resolution of our images. For example, C has a narrow range between ~ 2 and ~ 3 , and there is a lack of highly-concentrated values of C. Due to the low spatial resolution, the measurements of r_{80} and r_{20} are not robust. In particular, for highly-concentrated galaxies, the inner radius r_{20} is smaller than one pixel and is likely significantly overestimated, so C is underestimated. We will discuss these systematic effects using simulations below.

3.2.2. Systematic effects from simulations

In order to address how the low spatial resolution of the images affect the measurements of the parameters $CAGM_{20}$, or how reliable these parameters are (for $z \ge 6$ galaxies), we ran a series of slightly different simulations than we did in Section 3.1. Here we started with low-redshift real galaxies instead of model galaxies, because model galaxies are smooth, and do not cover a large range of the parameter space. For example, Ais zero for a noiseless single-component model galaxy. We chose to use the galaxy images in the library of galSVM (Huertas-Company et al. 2008, 2011). These low-redshift galaxies are large and bright, and their images have high S/N ratios. We randomly selected 1/10of the galaxies from this library. We then visually inspected these galaxies and removed those with possible foreground stars. Our final sample consists of 740 galaxies. Their morphological parameters are shown in Figure 6 (the scale in this figure is very different from that in Figure 5). These galaxies cover a large range of the parameter space of $CAGM_{20}$.

By rescaling these real galaxies in size and flux, we produced a large set of mock galaxies with high spatial resolution at $z \ge 6$, in the grid of magnitude $(-22 < M_{1500} < -19)$ and size $(0''.05 \le r_{\rm hl,in} \le 0''.4)$. The mock galaxies were then convolved with the PSF images and rebinned to match the pixel scales of our HST images. Finally, the rebinned mock galaxies were placed at many random positions in the blank regions of the HST images, and their morphological parameters were measured. Figure 7 shows an example of our simulation results (black dots). In this example, we preserved the relative magnitudes and sizes of the galaxies, and scaled the sample as a whole so that the scaled sample covers a similar range of magnitude and size as our $z \ge 6$ galaxies do. This is the best way to preserve the scatter of source distributions in the parameter space.

The comparison between Figure 6 and Figure 7 shows that the parameters of these mock galaxies measured in our low-resolution images are quite different from the intrinsic values: they occupy a smaller range of the parameter space. For example, the mock galaxies are significantly less concentrated in our images (2 < C < 3) in Figure 7 vs. 2 < C < 5 in Figure 6). As we already mentioned, C is significantly underestimated for

highly-concentrated galaxies because r_{20} is much smaller than one pixel in our HST images. Since M_{20} also describes galaxy concentration, it is underestimated as well (less negative here). This was also noticed by Lotz et al. (2006). Due to the low spatial resolution, our images cannot resolve subtle structures, so the mock galaxies show lower G coefficients, or more evenly distributed light. The limitation of G has been reported (e.g. Lisker 2008). Also because of the low resolution, the A values of the mock galaxies are much larger.

In Figure 7 we also plot our $z \ge 6$ galaxies (red circles). Their morphological measurements are directly taken from Figure 5. Their positions in the parameter space are quite consistent with those of the mock galaxies. This suggests that the $z \ge 6$ galaxies are possibly not intrinsically less concentrated, more asymmetric, or less even in light distribution. It is simply because our images do not meet one critical requirement to use these methods, i.e., large galaxy sizes compared to the PSF size, so our measurements have been systematically biased by the low-resolution images. On the other hand, these parameters are probably still meaningful for galaxies at similar redshifts, if they are measured in the HSTimages of the same depth and pixel size, i.e., for example, intrinsically more concentrated galaxies are still more concentrated in our HST images as measured by C. We demonstrate this using the simulations of the 740 mock galaxies shown in Figure 7. We first choose the galaxy pairs whose difference of the *measured* C (or any of $CASM_{20}$) are larger than 2σ (σ is the measurement uncertainty). Among these galaxy pairs, we further select the pairs in which the galaxy has larger *intrinsic* C than the other one still has larger *measured* C in the low-resolution images. The fraction of such pairs for C is 93%. The fractions for ASM_{20} are 70%, 98%, and 92%, respectively. If we increase 2σ to 3σ above, the fractions for $CASM_{20}$ increase to 97%, 78%, 100%, and 95%, respectively. This suggests that for the vast majority of galaxies, the low resolution of our HST images does not change their relative values of these morphological parameters. We emphasize that the *absolute* measured values of these parameters for $z \ge 6$ galaxies in HST images are not to be compared to the *absolute* measured parameters of lower-redshift galaxies.

3.2.3. Relations among the morphological parameters

The morphological parameters correlate with each other. We have seen such correlations in low-redshift galaxies (e.g. Lotz et al. 2004, 2006; Conselice & Arnold 2009). In Figure 6, the dashed lines are the best linear fits to the 740 low-redshift galaxies, and show the relations among $CAGM_{20}$. As expected, these parameters are correlated with each other. The moment M_{20} is correlated well with C, A, and G. Both M_{20} and C describe how the galaxy light is concentrated, and they are strongly correlated by definition. The relations between M_{20} and A & G reflect that more concentrated galaxies have more rotationally symmetric profiles and more unevenly distributed light. The Gini coefficient G is correlated with C and A, in addition to M_{20} . Its relations with C and A indicate that galaxies with more concentrated or more rotationally symmetric profiles tend to have more unevenly distributed light. These relations still exist among the 740 mock galaxies in Figure 7, since the relative values of the parameters are preserved as we discussed above. On the other hand, some relations become weaker with larger scatter. For example, the relations between C and AG are much flatter, mainly because the coverage of C has shifted from $2 \le C \le 5$ to $2 \le C \le 3$.

In Figure 5, the red dashed lines show the correlations of $CAGM_{20}$ for our $z \ge 6$ galaxies. These correlations are consistent with those for the mock galaxies in Figure 7. In particular, the relations between M_{20} and CAGare still fairly well preserved, and the relations between C and AG are as weak as shown in Figure 7. We will see in the next subsection that these corrections for relatively bright galaxies in our sample are better with smaller scatter (Figure 10). Therefore, although our measured $CAGM_{20}$ of the $z \ge 6$ galaxies are systematically biased by the low-resolution images, these parameters, when interpreted carefully, are still somewhat meaningful for galaxies measured in the HST images of the same depth and pixel size.

3.3. Interacting Systems

One of the interesting morphological topics is to study interacting/merging systems, which traces hierarchical mass assembly in the CDM scenario. A close visual inspection of Figure 1 shows that some galaxies are clearly extended with interacting or multi-component features. At lower redshift, these systems can be identified by pair counts (e.g., Le Fèvre et al. 2000), the CAS system (e.g., Conselice 2003), or the G and M_{20} parameters (e.g., Lotz et al. 2008). Galaxies at $z \ge 6$ are faint and small, so it is difficult to properly distinguish regular and interacting/merging systems. Although we have derived morphological parameters $CAGM_{20}$, they were biased, and their ability to identify merging systems at z > 6 has never been examined. Therefore, we identify interacting/merging systems by visual inspection. Visual classification was the earliest way for galaxy classification, and in many cases is still the best way to identify merging systems at high redshift.

We considered the following two types of galaxies as candidate interacting/merging systems: 1) galaxies with two or more distinct cores; and 2) galaxies with extended/elongated features and/or long tails. We identified these systems in 24 relatively bright ($M_{1500} \leq -20.5$ mag) galaxies. For fainter galaxies, our images are not deep enough to properly identify all faint components or extended features. Figures 8 and 9 show 12 galaxies that were identified as interacting systems with $M_{1500} \leq -20.5$ mag. Figure 8 shows 11 galaxies in the stacked (J+H) images, and Figure 9 shows the z = 6.96LAE (no. 62) in the two J bands. The red profiles are the surface brightness (SB) contours of the rest-frame UV emission. Six of them clear show double or multiple clumps, including no. 4, 24, 34, 49, 62, and 67. They usually have one bright core and one or more fainter clumps. No. 62 and 67 are particularly interesting. No. 62 in Figure 9 has almost two identical components, and no. 67 has three widely-separated cores. More discussion is given in the next subsection. Another 3 galaxies, including no. 36, 58, and 61, do not clearly show multiple clumps, but have long tails like tidal tails seen in low-redshift merging galaxies. The rest of the 12 systems (no. 15, 44, and 47) do not show multiple components or

tails, but they are rather extended and elongated. They could be in the end of the merging process.

We estimate the fraction of mergers among galaxies with $M_{1500} \leq -20.5$ mag. The fraction is 50%, or 38% if we exclude the three galaxies that do not show multiple components or tails. The fraction is even higher in the galaxies with $M_{1500} \leq -21$ mag. We have 18 galaxies in this magnitude range, and 10 of them are mergers. The fraction of mergers is 56%, or 39% if we exclude the three galaxies mentioned above. This is consistent with the fractions in the brightest galaxies at low redshift of $z \simeq 2 - 3$. For example, the fraction of mergers in the $M_B < -21$ mag galaxy sample of Conselice (2003) is 40%-50%. The typical fraction in the brightest galaxies in the sample of Law et al. (2012) is also $\sim 40\%$. The merger fraction in fainter galaxies is smaller, because merger systems consist of multiple components, and usually have stronger SFRs and UV emission. For the galaxies fainter than $M_{1500} \simeq -20.5$ mag in our sample, our images are deep enough to identify double-core systems with comparable emission. We find that these systems are very rare at $M_{1500} \ge -20.5$ mag. Studies in deeper fields have indicated a low merger fraction in low-luminosity galaxies. For example, Oesch et al. (2010) presented the morphology of 16 $z \ge 7$ LBGs in the HUDF, and only found two galaxies with extended features.

3.3.1. Notes on individual objects

Galaxy no. 62. No. 62 is a z = 6.964 LAE. It is the first spectroscopically-confirmed LAE at $z \sim 7$ (Iye et al. 2006). Figure 9 shows that it has two similar components in the both J bands. We used GALFIT to model the two components. Two Sérsic functions were fitted to the two components (two left-hand images in Figure 9) simultaneously. The middle images in Figure 9 show the best model fits, and the residuals are on the right-hand side. The two components in the both bands can be well described by the Sérsic function. The separation between the two cores is about 0'.'2 (~ 1 kpc) (see also Cai et al. 2011).

Galaxy no. 67. No. 67 is a LAE at z = 6.595. It was discovered as a giant LAE by Ouchi et al. (2009). It is one of the brightest galaxies in terms of both Ly α luminosity and UV continuum luminosity. The most striking feature is the three well separated cores lined up. The central core is relatively weak. This galaxy is clearly resolved in our ground-based z- and y-band images. The separate between the two side cores reaches $\simeq 1.2^{\circ}$, or 7 kpc. This is the largest separation we have seen at $z \ge 6$. This object also has strong emission in the IRAC bands, suggesting a large stellar mass of $\ge 10^{10} M_{\odot}$ (Ouchi et al. 2009).

3.3.2. Morphology of interacting systems

As we mentioned, the above interacting galaxies were identified with visual inspection. Morphological parameters were not used for selection because the criteria of interacting systems at $z \ge 6$ are unclear. We check our interacting systems in the parameter space in Figure 10. Figure 10 is similar to Figure 5, but only plots the galaxies with $M_{1500} \le -20.5$ mag in our sample. The red squares represent the interacting galaxies. Figure 10 shows that the interacting galaxies are almost indistinguishable from the rest in the parameter space. Asymmetry A is widely accepted as an efficient parameter to identify mergers at low redshift (e.g. Conselice 2003; Conselice & Arnold 2009; Lotz et al. 2006, 2008; Law et al. 2012). In Figure 10, however, A is not a good indicator of mergers any more, although the galaxies with the largest A in our sample are mostly interacting systems.

An interesting feature in Figure 10 is the better correlations among the morphological parameters compared to Figure 5. The scatter in the relations is smaller for these bright galaxies with $M_{1500} \leq -20.5$ mag. There are two explanations. One is that bright galaxies have higher S/N ratios and cover more pixels in the *HST* images, so their morphological measurements are more robust. The other one is that galaxies with different luminosities occupy slightly different parameter space, so a sample covering a smaller luminosity range shows a smaller scatter in Figure 10. The real reason is very likely the combination of the two. In any case, Figure 10 strengthens our earlier conclusion on the existence of strong correlations among $CAGM_{20}$ for our $z \geq 6$ galaxies.

4. $Ly\alpha$ MORPHOLOGY

In this section we will study the Ly α morphology of LAEs using our ground-based narrow-band images. We will not measure structure and morphology for individual galaxies, since they are mostly point-like objects in the ground-based images. Although these images have excellent PSF sizes of ~ 0".5 – 0".7, the PSF sizes are still much larger those of the HST images. Therefore, we will focus on Ly α halos around LAEs, which could extend many arcsec from the objects. We will also compare the positions of Ly α emission with those of UV continuum emission, and find any possible positional difference between the two.

4.1. $Ly\alpha$ Halos

Because of the resonant scattering of $Ly\alpha$ photons by neutral hydrogen, Ly α emission could form large diffuse $Ly\alpha$ halos around high-redshift galaxies. Steidel et al. (2011) first found very extended Ly α halos in a sample of luminous galaxies at 2 < z < 3. The galaxies were UV continuum-selected, but more than a half of them show net Ly α emission and ~ 20% have Ly α EW greater than 20 Å. They were able to find large $Ly\alpha$ emission halos (> 80 kpc) in the stacked images of all sub-samples of their galaxies. They further claimed that all LBGs would be classified as LAEs or $Ly\alpha$ blobs, if imaging data are deep enough to detect $Ly\alpha$ halos. Matsuda et al. (2012) confirmed the existence of extended $Ly\alpha$ halos around $z \simeq 3$ galaxies. They used more than 2000 LAEs at $z \simeq 3.1$, and grouped them into sub-samples based on luminosity and surface overdensity. They stacked narrowband $(Ly\alpha)$ images for each sub-sample, and found that all stacked images show extended (> 60 kpc) Ly α emission halos. Recently Feldmeier et al. (2013) found that the existence of $Ly\alpha$ halos around LAEs is not convincing. They also used a large sample of a few hundred LAEs at $z \simeq 2.1$ and 3.1. They paid particular attention to systematic effects from large-radius PSF and largescale flat fielding, etc. When these effects were taken into account, they did not find strong evidence of extended $Ly\alpha$ halos in the stacked narrow-band images at either

redshift. They tried a few ways to reconcile the discrepancy between their results and the previous results, yet the reason of the discrepancy is still not clear.

Stacking of narrow-band images has not been done for $z \geq 6$ galaxies. On the other hand, cosmological simulations have predicted the existence of extended Ly α emission around $z \geq 6$ galaxies (e.g., Zheng et al. 2011; Dijkstra & Kramer 2012; Jeeson-Daniel et al. 2012). For example, by including the resonant scattering of Ly α photons in both circumgalactic media and intergalactic medium (IGM), Zheng et al. (2011) showed that the Ly α emitting halo in a high-redshift galaxy can extend up to 1 Mpc. They further pointed out that such halos could be detected by stacking 100 $z \simeq 5.7$ LAEs in 4-hr exposure narrow-band images in the SXDS. Here we combine the narrow-band (Ly α) images of LAEs at $z \simeq 5.7$ and 6.5.

In order to detect diffuse $Ly\alpha$ halos around LAEs, we made use of all the known LAEs at $z \simeq 5.7$ and 6.5 in the SDF and SXDS fields from Kashikawa et al. (2011) and Ouchi et al. (2008, 2010). The narrow-band images were taken with the Subaru telescope. The total integration time in the NB816 and NB921 bands are 10 and 15 hr for the SDF (Kashikawa et al. 2004), and ~ 4 hr and ~ 10 hr (depth slightly varies among five pointings) for the SXDS (Ouchi et al. 2008, 2010). The data reduction of the SDF images were presented in Section 2 of Paper I. The reduction of the SXDS images were done in the same way. These narrow-band images have great depth with excellent PSF sizes of $0.5 \sim 0.7$. We rejected a small number of galaxies that are either very faint or blended with nearby bright objects. The final sample contains 43 LAEs at $z \simeq 5.7$ and 40 LAEs at $z \simeq 6.5$.

To stack these images, we first cut image stamps for all individual objects. We then re-sampled image stamps, so that the objects are in the centers of the images. After all other objects in the images were masked out, we co-added (average) the images with sigma-clipping (5 σ rejection). Figure 11 shows the stacked narrow-band $Ly\alpha$ images at the two redshifts. We did not subtract away continuum images. As we showed in Paper I, our LAEs have $Lv\alpha$ EWs greater than 20 Å. The median EW value is 80 Å, so their narrow-band photometry is completely dominated by the $Ly\alpha$ emission. Subtracting continuum images would significantly increase the background noise. In the above procedure we did not scale the objects to the same magnitude either, since we do not know how the SB of the Ly α halos scales with the total $Ly\alpha$ flux in a LAE. We performed a series of tests by stacking the images in different ways, including image stacking with median, with objects scaled to the same magnitude, and with bright objects only. The results are all very similar to the images shown in Figure 11. Our stacked images of LAEs reach a depth of roughly $1.2 \times 10^{-19} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2} (1\sigma)$ in both bands (see also Figure 12). We should point out that this depth is shallower than what we predicted assuming that the noise goes down with the square root of the image number. The reason is that the input images are not purely blank images, whose noise is dominated by background Poisson noise. These images are extremely deep and crowded, so LAEs have many nearby objects. Although we were able to mask out the bright pixels of

the nearby objects, we could not remove the light outside these pixels that may extend many pixels away. When we combine many images, these nearby objects contribute significant unavoidable noise to the co-added images.

Figure 11 clearly shows that the stacked $Ly\alpha$ emission is compact at the both redshifts, and do not show extended diffuse halos. As a reference, we also show in Figure 11 the stacked images of stars (point sources). The stars are chosen to be separated, bright, but not saturated, in the same narrow-band images. They are combined in the same way as we did for LAEs. Figure 12 shows the radial profiles of the stacked LAEs (solid lines with error bars) as well as the profiles of the stars (dotted lines). The Ly α profiles roughly follow the PSF profiles, and their SB reach zero at $r \ge 2''$. In the upper panel, the profile FWHM of the stacked $z \simeq 5.7$ LAE and star are $0''_{49}$ and $0''_{62}$, respectively. In the lower panel the two FWHM values are 0".61 for the $z \simeq 6.5$ LAE and $0\rlap.^{\prime\prime}\!77$ for the star. The LAE profiles are broader than the PSF sizes by $\sim 26\%$, and exhibit slightly longer tails than the PSF profiles do. This is simply because galaxies are not point sources, and indicates that the $Ly\alpha$ emission is resolved but not very extended. In the previous sections we show that our LAEs have a range of sizes in their rest-frame UV emission. Hathi et al. (2008) stacked broad-band HUDF images for $z \sim 4-6$ LBGs and found that the galaxy SB profiles are apparently broader than the PSF profiles. Finkelstein et al. (2011) found that two $z \simeq 4.4$ LAEs have large sizes in Ly α than in UV continuum. So the slightly broader $Ly\alpha$ SB profiles compared to the PSF are just the nature of galaxies, and are not likely caused by diffuse $Ly\alpha$ halos predicted or found in previous studies. The z = 5.7 profile seems to exhibit slightly more extended radius than its PSF profile. Given the 1σ limit of $\sim 1.2 \times 10^{-19} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$, its SB is still consistent with zero at r > 2''. Therefore, Figure 12 does not show convincing evidence of extended $Lv\alpha$ halos.

It is difficult to answer whether our stacked images are deep enough to detect $Ly\alpha$ halos if they do exist at the two redshifts. Our images are certainly deep enough to detect the LBG halos at $z \simeq 3$ reported by Steidel et al. (2011) and the LAE halos at $z \simeq 3.1$ in Matsuda et al. (2012). The depth to detect $z \ge 6$ LAE halos is observationally unknown. From cosmological simulations, Zheng et al. (2011) predicted Ly α halo sizes in z = 5.7LAEs. They found two characteristic scales for the halos. The inner steeper one extends to 3'' - 4'', and the outer flatter one extends to a few tens of arcsec. While our images are not deep enough to detect the outer halos, they are almost deep enough to detect the inner scale halos as seen in Figure 3 of Zheng et al. (2011), where the $Ly\alpha$ radial profiles are shown for LAEs in dark matter halos of ~ $10^{11} M_{\odot}$. As pointed out by Zheng et al. (2011), the size of diffuse $Ly\alpha$ emission also depends on the mass of dark matter halo. If the average mass of the dark matter halos in our LAEs is smaller than $10^{11} M_{\odot}$, our current data may not be able to detect the diffuse $Ly\alpha$ emission.

It is also likely that the Ly α halos (if they exist) have been diluted to a much lower level during the construction of the stacked images. The stacked images can properly recover Ly α halos only when halos are smoothly and symmetrically distributed around galaxies. If halos are highly asymmetric and/or clumpy, the emission of halos will be significantly diluted in average stacked images, and could totally disappear in median stacked images. From the observations of $z \sim 6$ quasars or cosmological simulations, we know that the distribution of IGM at $z \sim 6$ is inhomogeneous (e.g. Fan et al. 2006; Mesinger 2010). If the distribution of IGM affects the shape of Ly α halos (via resonant scattering), the distribution of Ly α halos is also likely asymmetric and clumpy.

Finally, it is possible that these LAEs do not have extended $Ly\alpha$ halos, or that their halo emission is not as strong as predicted by Zheng et al. (2011), especially when dust is taken into account (e.g. Finkelstein et al. 2011). Zheng et al. (2011) did not consider dust in their simulations. We know that high-redshift LAEs are not free of dust. In particular, the brightest galaxies may exhibit significant dust extinction, as implied by their UV colors (see Paper I). When $Ly\alpha$ photons are resonantly scattered by dusty neutral hydrogen, the Ly α emission is substantially reduced (Yajima et al. 2012). The reduction is more severe at larger distance from the object, because photons at larger distance need to pass through more dust before they escape. This process would significantly reduce the visibility of possible diffuse $Lv\alpha$ emission, and makes it much more difficult to detect it. A much larger sample of LAEs is needed to answer this question.

4.2. $Ly\alpha$ -Continuum Misalignment

The comparison between the positions of UV continuum and $Ly\alpha$ emission provides useful information on how Ly α photons escape from a galaxy. Ly α and UV continuum photons are usually come from the same starforming regions, although $Ly\alpha$ photons are likely more sensitive to the regions with more recent star-forming activity. As we mentioned earlier, $Ly\alpha$ emission is complicated by resonant scattering and IGM absorption. So the observed position of $Ly\alpha$ emission could be different from the position of UV continuum emission. For example, a large positional difference has been found in a z = 3.334 galaxy (Rauch et al. 2011). Due to the small sizes of high-redshift galaxies, current ground-based observations are not able to detect these positional differences. We rely on HST, which has observed a large number of high-redshift galaxies. HST observations were mostly made for rest-frame UV continuum emission, and there is usually no suitable HST narrow-band filters for $Ly\alpha$ emission. One example for the HST imaging of $Ly\alpha$ emission is the work of Finkelstein et al. (2011), who observed the Ly α emission of a small sample of $\simeq 4.4$ LAEs with a narrow-band filter. They did not find strong evidence of positional misalignment between UV and $Ly\alpha$ emission.

We used our large sample of LAEs to search for possible positional offsets between UV and Ly α emission at $z \simeq 5.7$ and 6.5. We used our HST images as UV continuum images and Subaru narrow-band images as Ly α images. As mentioned above, the narrow-band images have excellent PSF FWHM sizes around 0".5 ~ 0".7. During the construction of the HST images (Paper I), we have matched the coordinates of the HST images to those of the optical images. To avoid any large-scale variation, we refine the coordinates of the HST images. For each LAE, we found 10–20 nearby objects that are relatively bright and round. We then matched the positions of the nearby objects in the two sets of images. The typical refinement was smaller than the size of one pixel (0".06). The uncertainty in the object positions, derived from the distribution of the nearby objects, is about the size of 1– 2 pixels. We plot the Ly α positions on top of the UV continuum positions. Figure 13 shows a few examples of bright galaxies. The red profiles are the contours of the UV emission SB, and the green crosses indicate the positions (and 1 σ uncertainties) of the Ly α emission.

For the majority of the LAEs in our sample, the positions of the UV continuum and $Ly\alpha$ emission agree with each other. In particular, we do not find positional misalignment at a significance level of $> 2\sigma$ among almost all compact and round LAEs. Object no. 3 in Figure 13 is a typical example, in which the center of the $Lv\alpha$ emission is close to the center of the continuum emission. For the merging/interacting systems, however, we see significant positional differences. Figure 13 shows the examples of these systems. They exhibit a variety of $Ly\alpha$ positions relative to the peak positions¹² of the UV continuum emission: 1) Ly α positions close to the positions of the brightest components in the UV images, including galaxies no. 15, 47, and 61; 2) Ly α positions close to the positions of the fainter components or merger tails in the UV images, including no. 14, 44, and 58; 3) Ly α positions are somewhere between the positions of the bright and faint components, but closer to the bright components, including no. 4 and 49. No. 62 and 67 are again interesting. The Ly α position of no. 62 is roughly in the middle of the two similarly bright components. No. 67 does not show three distinct $Ly\alpha$ emission clumps, as its UV emission does. Instead, it shows a single bright $Ly\alpha$ emission core with some extended features. The $Ly\alpha$ center is not at any of the three UV clumps. It is between the left and central clumps, and slightly closer to the central one.

Our results suggest that in compact galaxies the observed location of $Ly\alpha$ emission does not deviate from its original position, while in merging/interacting systems, the observed $Ly\alpha$ location could be significantly different from its original position, without preferential positions of offsets. If the final location of $Ly\alpha$ emission is determined by the process of resonant scattering, our results can be explained by the interstellar medium (ISM) distribution (e.g. Finkelstein et al. 2011). In a non-disturbed galaxy, the ISM distribution is relatively symmetric around the object (not necessarily uniform, it could be clumpy). The random scattering of photons does not have preferred directions, so the observed location is still close to its original location. In an interacting system, the ISM is re-distributed by merging activity. The distribution of the disturbed ISM is therefore not symmetric any more. This results in an offset of the observed Ly α position.

This result could have an impact on spectroscopy of bright LBGs at $z \ge 6$ (e.g. Stark et al. 2011; Curtis-Lake et al. 2012). While most of the positional offsets in our sample are smaller than 0'.2, at least two

 $^{^{12}}$ For an interacting galaxy, the position of its peak emission (as seen from the SB contour) could be very different from the position of the overall galaxy emission measured by, for example, fitting Gaussian to the marginal x,y distributions (used by IRAF DAOPHOT). Here our positions refer to the positions of peak emission.

are around $0''_{...3} - 0''_{...4}$. If one uses $a \leq 1''$ slit to identify their Ly α emission lines, based on the positions of continuum emission, one could miss them due to the large offsets. However, glaxies with such large offsets are very rare, so this result will not largely reduce the success rate of identifying Ly α lines in LBGs.

5. DISCUSSION AND SUMMARY

The comparison in Section 3.1 shows that the sizes of our galaxies are consistent with those of bright $z \geq 6$ LAEs and LBGs in previous studies. This disagrees with the claim of Dow-Hygelund et al. (2007) that LAEs are more compact than LBGs. This is likely because the Dow-Hygelund et al. (2007) sample is very small while LAEs have a large range of sizes. It is indeed difficult to make proper comparisons without a large sample, as the galaxy size depends on redshift and luminosity. Most of our galaxies are LAEs. Another large LAE sample is the $z \simeq 5.7$ LAE sample by Taniguchi et al. (2009), who observed a number of LAEs at $z \simeq 5.7$ with the HST ACS F814W filter. Although this filter includes $Ly\alpha$ emission, it is so wide that its emission is dominated by UV continuum. They found that the average size of the sample is 0.15, similar to ours. Among photometrically-selected LBG samples at $z \ge 6$, a recent large sample is the sample of Grazian et al. (2012). This sample contains a number of bright (as well as faint) LBGs at $z \sim 6-7$ in the CANDELS field. The sizes of these galaxies are well consistent with those of our galaxies at the same luminosities. Therefore, we conclude that LAEs and LBGs with similar luminosities have similar physical sizes.

In this paper (and Paper I), galaxies found by the narrow-band technique are defined as LAEs and those found by the dropout technique are defined as LBGs. As we discussed in Paper I, this classification only reflects the methodology that we apply to select galaxies. It does not mean that the two types of galaxies are intrinsically different. In Section 5.3 of Paper I, when we derived the UV continuum luminosity function of LAEs, we used another popular definition of LAEs based on the $Ly\alpha$ EW, i.e., a galaxy is a LAE if its $Ly\alpha$ EW is greater than 20 Å. With this definition, almost all the galaxies in our sample are LAEs. This definition is physically more meaningful, but observationally difficult, because one can easily obtain a flux-limited sample, not a EW-limited sample. We have 16 LBGs (former definition) in our sample. They are not typical LBGs. They are spectroscopically confirmed, and thus only represent those with strong $Ly\alpha$ emission. In Paper I we found that our LAEs and LBGs are indistinguishable in many aspects of the Ly α and UV continuum properties. In this paper, we further found that these LAEs and LBGs have similar physical sizes and morphological parameters. This confirms one of our conclusions in Paper I that LAEs are a subset of LBGs with strong $Ly\alpha$ emission lines.

This paper is the second in a series presenting the physical properties of a large sample of spectroscopicallyconfirmed galaxies at $z \ge 6$. The sample consists of 51 LAEs and 16 LBGs, and represents the most luminous galaxies in terms of Ly α luminosity (for LAEs) or UV continuum luminosity (for LBGs) in this redshift range. In Paper I we derived various properties of rest-frame UV continuum and Ly α emission. In this paper we have conducted a detailed structural and morphological study of the galaxies using deep HST near-IR images and Subaru narrow-band images. In order to measure the morphology of rest-frame UV continuum emission, we constructed a stacked HST image for each object by combining its *J*- and *H*-band images. UV morphology was then measured for those with $> 10\sigma$ detections (roughly corresponding to $M_{1500} \leq -19.5$ mag) in the stacked HST images.

We have used half-light radius to describe the sizes of galaxies. The intrinsic sizes $r_{\rm hl,in}$ (at $M_{1500} \leq -19.5$ mag) in our sample, after correction for PSF broadening, are from $\leq 0'.05~(\leq 0.3~{\rm kpc})$ to $\sim 0''.3~(\sim 1.7~{\rm kpc})$, with a median value of $0'.16~(\sim 0.9~{\rm kpc})$. These values are consistent with those of bright, photometrically-selected LBGs at similar redshifts. Additionally, more luminous galaxies tend to have larger sizes, exhibiting a weak size-luminosity relation $r_{\rm hl,in} \propto L^{0.14}$. The slope 0.14 is significantly flatter than those in fainter LBG samples. Our objects show a wide range of morphology in the HST images, including compact galaxies and double/multiple component systems. The brightest galaxies in the sample have a large fraction of merging/interacting systems. The fraction of mergers reaches 40 - 50% at $M_{1500} \leq -20.5$ mag.

We have tried to describe the structure and morphology of our $z \ge 6$ galaxies using nonparametric methods, including the CAS system, the Gini and M_{20} parameters. Compared to low-redshift galaxies, these galaxies occupy a smaller range in the parameter space of $CAGM_{20}$. Our simulations show that the measurements of these parameters have been systematically biased due to the low resolution of the HST images. On the other hand, the relative values of these morphological parameters are likely preserved for the vast majority of galaxies. In addition, we found all expected correlations among $CAGM_{20}$ for the $z \ge 6$ galaxies. These suggest that the parameters probably still meaningful for galaxies at similar redshifts, if they are measured in the HST images of the same depth and pixel size.

We, for the first time, searched for $Ly\alpha$ emission halos around z > 6 galaxies in narrow-band images. We combined a large number of narrow-band images for LAEs at $z \simeq 5.7$ and 6.5, respectively. The stacked images reached a depth of $\sim 1.2 \times 10^{-19} \rm \ erg \ s^{-1} \ cm^{-2} \ arcsec^{-2}$ (1σ) . We did not find evidence of extended diffuse Ly α emission as predicted by cosmological simulations. It is possible that our images are still not deep enough to detect $Ly\alpha$ emission halos, or that the $Ly\alpha$ halo emission has been diluted to a much lower level during the construction of the stacked images. It is also possible that the halo emission is not as strong as predicted. A much larger LAE sample is needed to solve this question. We also investigated positional differences between the rest-frame UV continuum emission and $Ly\alpha$ emission in LAEs, using the HST images and optical narrow-band images. While in compact LAEs the two positions are well consistent, in some merging galaxies show significant positional differences, with no preferred directions of offsets. It was explained by the distribution of the disturbed ISM.

Support for this work was provided by NASA through Hubble Fellowship grant HST-HF-51291.01 awarded by

the Space Telescope Science Institute (STScI), which is operated by the Association of Universities for Research in Astronomy, Inc., for NASA, under contract NAS 5-26555. L.J., E.E., M.M, and S.C. also acknowledge the support from NASA through awards issued by STScI (HST PID: 11149, 12329, 12616) and by JPL/Caltech (Spitzer PID: 40026,70094).

Facilities: HST (NICMOS, WFC3), Spitzer (IRAC), Subaru (Suprime-Cam)

REFERENCES

- Abraham, R. G., Tanvir, N. R., Santiago, B. X., et al. 1996, MNRAS, 279, L47
- Abraham, R. G., & van den Bergh, S. 2001, Science, 293, 1273 Abraham, R. G., van den Bergh, S., & Nair, P. 2003, ApJ, 588, 218
- Bershady, M. A., Jangren, A., & Conselice, C. J. 2000, AJ, 119, 2645
- Bertin, E., & Arnouts, S. 1996, A&AS, 117, 393
- Blanton, M. R., & Moustakas, J. 2009, ARA&A, 47, 159
- Bouwens, R. J., Illingworth, G. D., Blakeslee, J. P., & Franx, M. 2006, ApJ, 653, 53
- Brinchmann, J., Abraham, R., Schade, D., et al. 1998, ApJ, 499, 112
- Cai, Z., Fan, X., Jiang, L., et al. 2011, ApJ, 736, L28
- Carlberg, R. G., Cohen, J. G., Patton, D. R., et al. 2000, ApJ, 532. L1
- Cassata, P., Giavalisco, M., Guo, Y., et al. 2010, ApJ, 714, L79
- Cibinel, A., Carollo, C. M., Lilly, S. J., et al. 2012, arXiv:1206.6108
- Cole, S., Lacey, C. G., Baugh, C. M., & Frenk, C. S. 2000, MNRAS, 319, 168
- Conselice, C. J., 2003, ApJS, 147, 1 Conselice, C. J., & Arnold, J. 2009, MNRAS, 397, 208
- Cooke, J., Berrier, J. C., Barton, E. J., Bullock, J. S., & Wolfe, A. M. 2010, MNRAS, 403, 1020
- Cowie, L. L., Hu, E. M., & Songaila, A. 2011, ApJ, 735, L38
- Curtis-Lake, E., McLure, R. J., Pearce, H. J., et al. 2012, MNRAS, 422, 1425
- Dijkstra, M., & Kramer, R. 2012, MNRAS, 424, 1672
- Dow-Hygelund, C. C., Holden, B. P., Bouwens, R. J., et al. 2007, ApJ, 660, 47
- Driver, S. P., Windhorst, R. A., Ostrander, E. J., et al. 1995, ApJ, 449, L23
- Driver, S. P., Fernandez-Soto, A., Couch, W. J., et al. 1998, ApJ, 496, L93
- Driver, S. P., Liske, J., Cross, N. J. G., De Propris, R., & Allen, P. D. 2005, MNRAS, 360, 81
- Fan, X., Carilli, C. L., & Keating, B. 2006, ARA&A, 44, 415
- Feldmeier, J., Hagen, A., Ciardullo, R., et al. 2013, arXiv:1301.0462
- Ferguson, H. C., Dickinson, M., Giavalisco, M., et al. 2004, ApJ, 600, L107
- Finkelstein, S. L., Cohen, S. H., Windhorst, R. A., et al. 2011, ApJ. 735. 5
- Furusawa, H., Kosugi, G., Akiyama, M., et al. 2008, ApJS, 176, 1
- Giavalisco, M., Steidel, C. C., & Macchetto, F. D. 1996, ApJ, 470, 189
- Giavalisco, M., Ferguson, H. C., Koekemoer, A. M., et al. 2004, ApJ, 600, L93
- Grazian, A., Castellano, M., Fontana, A., et al. 2012, A&A, 547, A51
- Grogin, N. A., Kocevski, D. D., Faber, S. M., et al. 2011, ApJS, 197.35
- Gronwall, C., Bond, N. A., Ciardullo, R., et al. ApJ, 743, 9
- Hathi, N. P., Jansen, R. A., Windhorst, R. A., et al. 2008, AJ, 135, 156
- Häussler, B., McIntosh, D. H., Barden, M., et al. 2007, ApJS, 172, 615
- Huertas-Company, M., Rouan, D., Tasca, L., Soucail, G., & Le Fèvre, O. 2008, A&A, 478, 971
- Huertas-Company, M., Aguerri, J. A. L., Bernardi, M., Mei, S., & Sánchez Almeida, J. 2011, A&A, 525, A157
- Iye, M., Ota, K., Kashikawa, N., et al. 2006, Nature, 443, 186
- Jeeson-Daniel, A., Ciardi, B., Maio, U., et al. 2012, MNRAS, 424, 2193
- Jiang, L., Egami, E., Kashikawa, N., et al. 2011, ApJ, 743, 65
- Jiang, L., Egami, E., Mechtley, M., et al. 2013, ApJ, in press (arXiv:1303.0024)

- Kashikawa, N., Shimasaku, K., Yasuda, N., et al. 2004, PASJ, 56, 1011
- Kashikawa, N., Shimasaku, K., Malkan, M. A., et al. 2006, ApJ, 648, 7
- Kashikawa, N., Shimasaku, K., Matsuda, Y., et al. 2011, ApJ, 734, 119
- Kashikawa, N., Nagao, T., Toshikawa, J., et al. 2012, ApJ, 761, 85 Koekemoer, A. M., Faber, S. M., Ferguson, H. C., et al. 2011,
- ApJS, 197, 36 Komatsu, E., Smith, K. M., Dunkley, J., et al. 2011, ApJS, 192, 18
- Law, D. R., Steidel, C. C., Erb, D. K., et al. 2007, ApJ, 656, 1
- Law, D. R., Steidel, C. C., Shapley, A. E., et al. 2012, ApJ, 745, 85
- Le Fèvre, O., Abraham, R., Lilly, S. J., et al. 2000, MNRAS, 311, 565
- Lilly, S., Schade, D., Ellis, R., et al. 1998, ApJ, 500, 75
- Lisker, T. 2008, ApJS, 179, 319
- Lotz, J. M., Primack, J., & Madau, P. 2004, AJ, 128, 163
- Lotz, J. M., Madau, P., Giavalisco, M., Primack, J., & Ferguson, H. C. 2006, ApJ, 636, 592
- Lotz, J. M., Davis, M., Faber, S. M., et al. 2008, ApJ, 672, 177 Lowenthal, J. D., Koo, D. C., Guzman, R., et al. 1997, ApJ, 481, 673
- Malhotra, S., Rhoads, J. E., Finkelstein, S. L., et al. 2012, ApJ, 750, L36
- Matsuda, Y., Yamada, T., Hayashino, T., et al. 2012, MNRAS, 425, 878
- Mesinger, A. 2010, MNRAS, 407, 1328
- Mosleh, M., Williams, R. J., Franx, M., et al. 2012, ApJ, 756, L12
- Nagao, T., Taniguchi, Y., Kashikawa, N., et al. 2004, ApJ, 613, L9
- Nagao, T., Kashikawa, N., Malkan, M. A., et al. 2005, ApJ, 634, 142
- Nagao, T., Murayama, T., Maiolino, R., et al. 2007, A&A, 468, 877
- Oesch, P. A., Bouwens, R. J., Carollo, C. M., et al. 2010, ApJ, 709, L21
- Oke, J. B., & Gunn, J. E. 1983, ApJ, 266, 713
- Ono, Y., Ouchi, M., Curtis-Lake, E., et al. 2012, arXiv:1212.3869
- Ouchi, M., et al. 2008, ApJS, 176, 301
- Ouchi, M., Mobasher, B., Shimasaku, K., et al. 2009, ApJ, 706, 1136
- Ouchi, M., Shimasaku, K., Furusawa, H., et al. 2010, ApJ, 723, 869
- Ota, K., Kashikawa, N., Malkan, M. A., et al. 2008, arXiv:0804.3448
- Peng, C. Y., Ho, L. C., Impey, C. D., & Rix, H.-W. 2002, AJ, 124.266
- Pirzkal, N., Malhotra, S., Rhoads, J. E., & Xu, C. 2007, ApJ, 667, 49
- Rauch, M., Becker, G. D., Haehnelt, M. G., et al. 2011, MNRAS, 418, 1115
- Ravindranath, S., Giavalisco, M., Ferguson, H. C., et al. 2006, ApJ, 652, 963
- Schade, D., et al. 1999, ApJ, 525, 31
- Shi, Y., Rieke, G., Lotz, J., & Perez-Gonzalez, P. G. 2009, ApJ, 697, 1764
- Shimasaku, K., Kashikawa, N., Doi, M., et al. 2006, PASJ, 58, 313
- Stanway, E. R., Glazebrook, K., Bunker, A. J., et al. 2004, ApJ, 604. L13
- Stark, D. P., Ellis, R. S., & Ouchi, M. 2011, ApJ, 728, L2
- Steidel, C. C., Bogosavljević, M., Shapley, A. E., et al. 2011, ApJ, 736.160
- Taniguchi, Y., Ajiki, M., Nagao, T., et al. 2005, PASJ, 57, 165
- Taniguchi, Y., Murayama, T., Scoville, N. Z., et al. 2009, ApJ, 701, 915
- Taylor-Mager, V. A., Conselice, C. J., Windhorst, R. A., & Jansen, R. A. 2007, ApJ, 659, 162

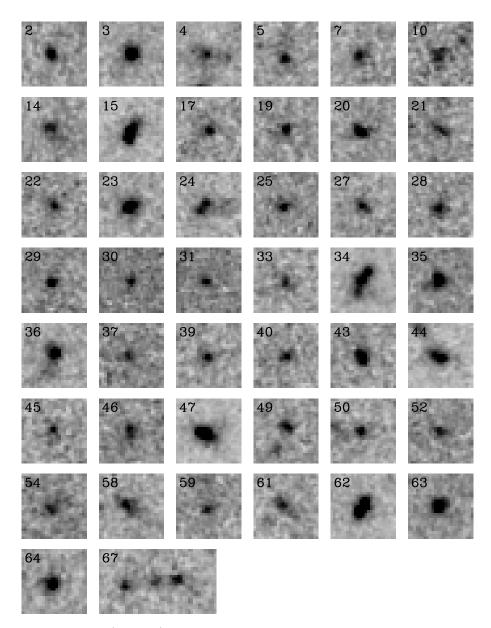


FIG. 1.— Thumbnail images of the 44 (out of 67) galaxies that have more than 10σ detections in the stacked *HST* images, shown in order of increasing redshift. The numbers of the galaxies plotted in this figure correspond to the galaxy numbers given in Table 1 of Paper I. The size of the first 43 images is 25 pixel square, which corresponds to 2.5×2.5 in NICMOS images (no. 25, 30, and 31), or 1.5×1.5 in WFC3 images (the rest). The size of no. 67 is 45 by 25 pixels (or 2.7×1.5).

Toshikawa, J., Kashikawa, N., Ota, K., et al. 2012, ApJ, 750, 137
van den Bergh, S., Cohen, J. G., Hogg, D. W., & Blandford, R. 2000, AJ, 120, 2190

van der Wel, A., Bell, E. F., Häussler, B., et al. 2012, ApJS, 203, 24

- Venemans, B. P., Röttgering, H. J. A., Miley, G. K., et al. 2005, A&A, 431, 793
- White, S. D. M., & Rees, M. J. 1978, MNRAS, 183, 341

Windhorst, R. A., Taylor, V. A., Jansen, R. A., et al. 2002, ApJS, 143, 113

- Windhorst, R. A., Hathi, N. P., Cohen, S. H., et al. 2008, Advances in Space Research, 41, 1965
- Windhorst, R. A., Cohen, S. H., Hathi, N. P., et al. 2011, ApJS, 193, 27
- Yajima, H., Li, Y., Zhu, Q., & Abel, T. 2012, MNRAS, 424, 884
- Zheng, Z., Cen, R., Weinberg, D., Trac, H., & Miralda-Escudé, J. 2011, ApJ, 739, 62

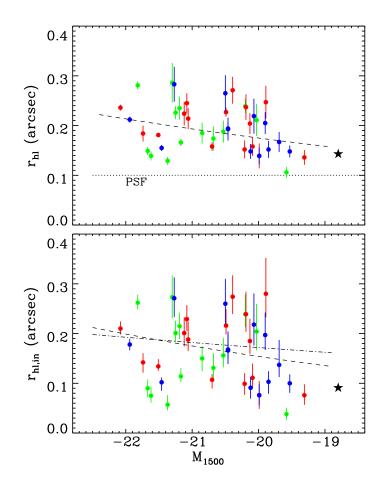


FIG. 2.— Half-light radius as a function UV continuum luminosity M_{1500} . The blue and red circles represent the LAEs at $z \simeq 5.7$ and 6.5 (including $z \simeq 7$), respectively, and the green circles represent the LBGs at $z \simeq 6$. The black star represents the stacked object of the faint galaxies, and is arbitrarily put at $M_{1500} = -18.8$ mag. The upper panel shows the measured radius r_{hl} without correction for PSF broadening. The error bars reflect the measurement uncertainties derived from simulations in Section 3.1. The dotted line indicates the PSF size in our *HST* WFC3 images. The dashed line (best log-linear fit) illustrates the weak relation between r_{hl} and M_{1500} ($r_{hl} \propto L^{0.11\pm0.02}$). The lower panel shows the intrinsic radius $r_{hl,in}$ after correction for PSF broadening with simulations in Section 3.1. The error bars include both measurement and systematic uncertainties. The dashed line (best log-linear fit to all data points) shows the weak relation between $r_{hl,in} \propto L^{0.14\pm0.03}$) for the whole sample. The dash-dotted line (best log-linear fit to the data points at $M_{1500} \leq -20$ mag) suggests little correlation ($r_{hl,in} \propto L^{0.06\pm0.03}$) between size and luminosity in luminous galaxies.

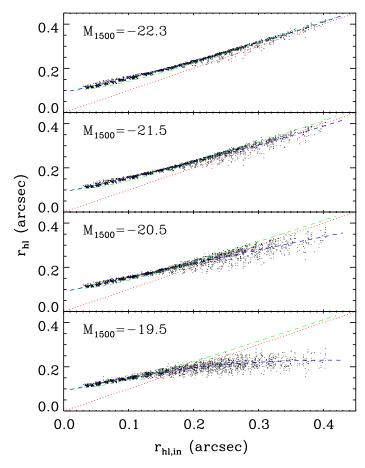


FIG. 3.— Measured half-light radius $r_{\rm hl}$ as a function of intrinsic half-light radius $r_{\rm hl,in}$ at four different magnitudes, derived from our simulations in Section 3.1. The red dotted lines indicate the equality relation. At small sizes, $r_{\rm hl}$ is significantly larger than $r_{\rm hl,in}$ due to PSF broadening. At large sizes $(r_{\rm hl,in} \ge 0.2^{\prime\prime} - 0.2^{\prime\prime})$, $r_{\rm hl}$ starts to fall short of $r_{\rm hl,in}$. This happens at smaller sizes for fainter galaxies. We illustrate this trend by displaying the best second-order polynomial fit to the data points (blue dashed lines). The green dashed lines show the relation of $r_{\rm hl,in}^2 = r_{\rm hl}^2 - r_{\rm PSF}^2$, which is a good approximation at small sizes and/or high luminosities, but underestimates $r_{\rm hl,in}$ elsewhere.

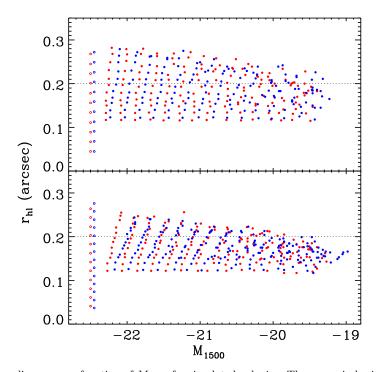


FIG. 4.— Measured half-light radius r_{h1} as a function of M_{1500} for simulated galaxies. The open circles indicate the input intrinsic sizes $r_{h1,in}$ of the galaxies in our simulations. The filled circles are the measured sizes r_{h1} at different luminosities. The two colors red and blue indicate two different HST images that the mock galaxies are placed in. The dotted lines are used to guide the eye. The two panels are for two sets of Sérsic index n: 1 and 2 times the measured n from our $z \ge 6$ galaxies. They show that fainter galaxies (with the same intrinsic size) appear to be smaller, as also shown in Figure 3. For the same reason, larger galaxies (with the same intrinsic luminosity) appear to be slightly fainter (see discussion in Section 3.1).

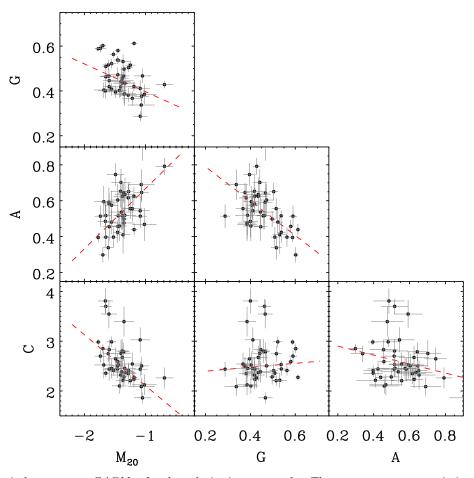


FIG. 5.— Morphological parameters $CAGM_{20}$ for the galaxies in our sample. The measurement uncertainties were estimated from simulations in Section 3.2.1. Compared to low-redshift galaxies, our galaxies occupy a much narrower range in the parameter space due to the low resolution of the images. Nevertheless, each parameter is correlated with one or more of the other parameters. The red dashed lines are the best linear fits to the relations.

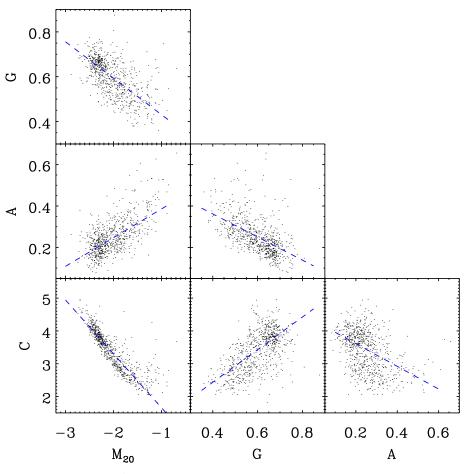


FIG. 6.— Morphological parameters $CAGM_{20}$ for 740 low-redshift bright galaxies selected from the library of galSVM (Huertas-Company et al. 2008, 2011). These galaxies cover a large range of the parameter space. The blue dashed lines are the best linear fits to the data points, showing the correlations among these parameters. Note that the scale in this figure is very different from that in Figure 5.



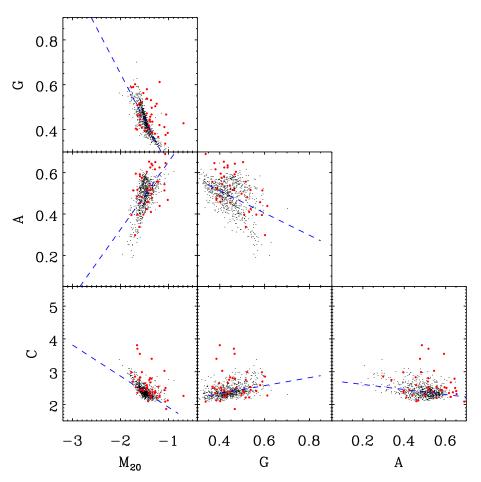


FIG. 7.— Morphological parameters $CAGM_{20}$ of 740 mock galaxies at $z \ge 6$ (black dots) measured in our HST images. The red circles represent our $z \ge 6$ galaxies (same as Figure 5). The scale in this figure is the same as the scale in Figure 6. The mock galaxies were produced from the 740 low-redshift galaxies (Figure 6) using simulations in Section 3.2.2. Compared to Figure 6, the mock galaxies occupy a smaller range of the parameter space due to the low resolution of our HST images. The blue dashed lines are the best linear fits to the black dots, showing the correlations among these parameters. The positions in the parameter space are consistent between the $z \ge 6$ and mock galaxies, suggesting that the $z \ge 6$ galaxies are probably not intrinsically less concentrated, nor more asymmetric.

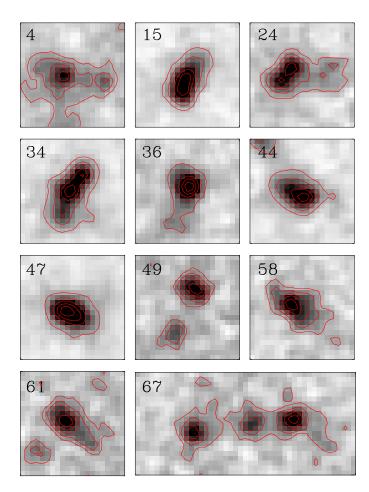


FIG. 8.— Interacting/merging galaxies identified in our HST images. The image size is $1.26 \times 1.26 \times 1.26$. The red profiles are the contours of SB. Each contour starts at 85% of the peak value with an interval of 20%. These galaxies have extended and elongated features, and/or have double or multiple components.

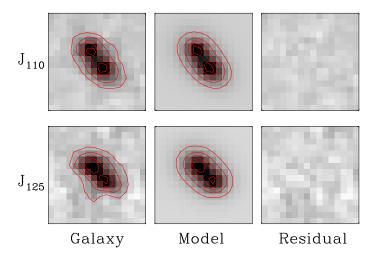


FIG. 9.— Images of the z = 6.96 LAE (no. 62). The two left-most images show the galaxy in the two J bands (J_{110} and J_{125}). The two images in the middle are the best-fit model (two Sérsic functions for the two components) galaxies. The residuals are on the right-hand side. The separation between the two components is about 0''.2 (~ 1 kpc).

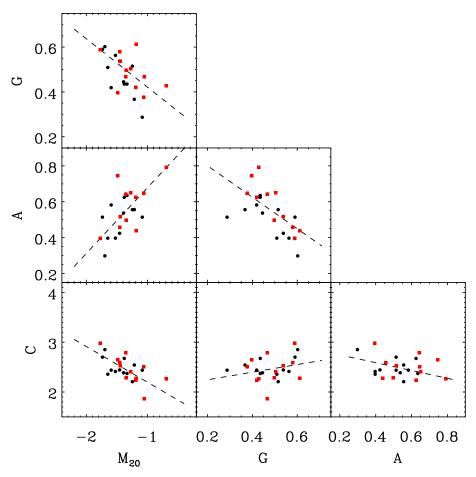


FIG. 10.— Same as Figure 5, but for galaxies with $M_{1500} \leq -20.5$ mag. The red squares represent the interacting galaxies identified in our sample. The figure also shows better corrections (smaller scatter) among the parameters compared to Figure 5.

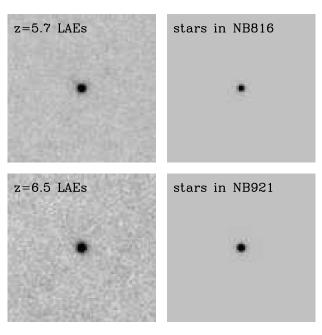


FIG. 11.— Stacked narrow-band images for LAEs and stars (or point sources) in two bands NB816 and NB921. The image size is $20'' \times 20''$. The images have the same intensity scale. The PSF FWHM sizes derived from the two stacked stars are 0''49 and 0''61, respectively. The FWHM of the stacked LAEs are 0''61 and 0''77, respectively. They are larger than the PSF sizes by 26%, because LAEs are not point sources. The stacked LAEs do not show diffuse Ly α halos.

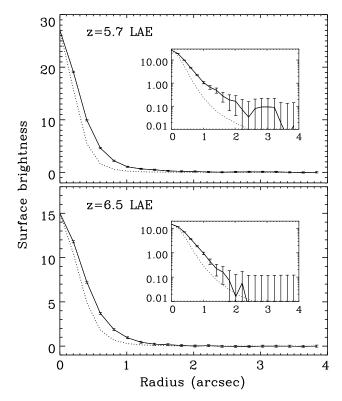


FIG. 12.— Radial profiles of the stacked images in Figure 11. The units of SB is $10^{-19} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$. The solid profiles with 1σ error bars represent the stacked LAEs, and the dashed profiles represent the stacked stars. The insets show the radial profiles on a log scale. The LAE profiles are broader than the PSF sizes, and exhibit slightly longer tails than the PSF profiles do, meaning that the Ly α emission is resolved. This is because galaxies are not point sources, as explained in Section 4.1. At $r \geq 2''$, the SB in both cases is consistent with zero within 1σ errors. We do not see very extended halos of Ly α emission.



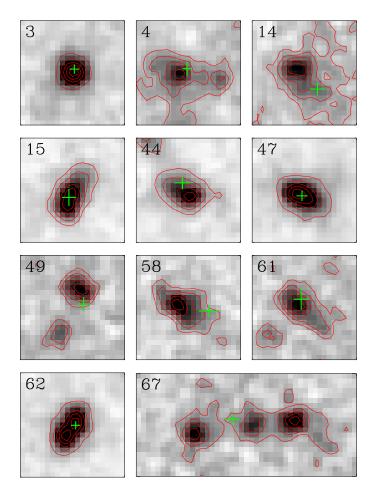


FIG. 13.— Positions of UV continuum and Ly α emission for a sample of bright galaxies. The red contours and their scales are similar to those shown in Figure 8. They display the UV continuum emission seen in the HST images. The green crosses indicate the positions (and 1σ uncertainties) of the Ly α emission from our ground-based narrow-band images. The first object no. 3 represents a typical compact galaxy, whose positions of UV continuum and Ly α emission agree with each other. The rest of the objects are merging/interacting systems, which show a variety of Ly α position offsets relative to the UV continuum positions, including significant positional misalignment.